

1 Cirrus Cloud Seeding has Potential to Cool Climate

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2 Cirrus clouds, thin ice clouds in the upper troposphere, have a net warm-
3 ing effect on Earth's climate. Consequently, a reduction in cirrus cloud amount
4 or optical thickness would cool the climate. Recent research indicates that
5 by seeding cirrus clouds with particles that promote ice nucleation, their life-
6 times and coverage could be reduced. We have tested this hypothesis in a
7 global climate model with a state-of-the-art representation of cirrus clouds,
8 and find that cirrus cloud seeding has the potential to cancel the entire warm-
9 ing caused by human activity from pre-industrial times to present day. How-
10 ever, the desired effect is only obtained for seeding particle concentrations
11 that lie within an optimal range. With lower than optimal particle concen-
12 trations a seeding exercise would have no effect. Moreover, a higher than op-
13 timal concentration results in an over-seeding that could have the deleteri-
14 ous effect of prolonging cirrus lifetime and contributing to global warming.

1. Introduction

15 With the realization that Earth's climate is changing at a rapid pace, a number of
16 mechanisms through which climate could artificially be stabilized have been proposed
17 in the literature. Climate sensitivity, defined as the equilibrium surface temperature re-
18 sponse to a doubling of atmospheric CO₂, is poorly constrained, and very high climate
19 sensitivities cannot currently be ruled out [*Roe and Baker, 2007*]. This, combined with
20 what seems to be a difficult prospect of curbing anthropogenic CO₂ emissions [*Davis*
21 *et al.*, 2010], the main cause of modern climate change, has led many to propose cli-
22 mate engineering as a cooling mechanism [*Keith, 2001; Boyd, 2008*]. Carbon capturing
23 and sequestration is one example of climate engineering that would directly target the
24 problem of rising atmospheric CO₂ concentrations [*Metz et al.*, 2005]. Another class of
25 climate engineering proposals is often termed solar radiation management (SRM), be-
26 cause rather than reducing Earth's greenhouse effect, their purpose is to increase Earth's
27 albedo/reflectivity. SRM strategies include stratospheric sulphur injection, mimicking vol-
28 canic eruptions [*Crutzen, 2006*], and enhancement of marine stratocumulus cloud albedo
29 via sea salt injection [*Latham, 1990*]. Both mechanisms have been the focus of many
30 recent studies [*Rasch et al.*, 2008; *Wang et al.*, 2011], and several complications have been
31 identified. Examples are changes to the local and regional hydrological cycles [*Ricke et*
32 *al.*, 2010], as well as stratospheric ozone depletion in the case of stratospheric sulphur
33 injection [*Tilmes et al.*, 2008]. Here we address a climate engineering mechanism that has
34 so far not been tested; the perturbation of cirrus clouds to reduce their lifetime and op-
35 tical thickness, thereby cooling Earth's climate. This idea was first put forth by *Mitchell*

36 *and Finnegan* [2009], and builds on the fact that spontaneous freezing of liquid solution
37 droplets requires high water vapor partial pressures that well exceed that of saturation
38 with respect to a plane ice surface (i.e. supersaturation, S_i). Spontaneous freezing of
39 droplets is a stochastic process that is referred to as homogeneous nucleation [*Koop et*
40 *al.*, 2000]. The homogeneous nucleation rate decreases with increasing temperature (T),
41 and for T higher than about -35°C , homogeneous ice nucleation does not occur in the
42 atmosphere [*Pruppacher and Klett*, 1997]. The presence of a substrate to facilitate the
43 formation of tiny ice crystals can significantly lower the supersaturation required for ice
44 formation, a process known as heterogeneous ice nucleation. Certain insoluble particles
45 can provide such substrates in the atmosphere, and are termed ice nuclei (IN). Examples
46 of natural IN are mineral dust particles, as well as certain primary biological particles
47 [*Pruppacher and Klett*, 1997]. Bismuth tri-iodide (BiI_3) is an example of an artificial IN,
48 and has been suggested as cirrus seeding material [*Mitchell and Finnegan*, 2009]. It has
49 been suggested that BiI_3 can initiate freezing at a supersaturation as low as 5%, while
50 homogeneous ice nucleation requires a supersaturation of the order of 50% at typical cirrus
51 temperatures. Due to the low concentration of IN in the upper troposphere (UT), homo-
52 geneous freezing is thought to dominate cirrus cloud formation [*Kaercher and Lohmann*,
53 2003; *Mitchell et al.*, 2011]. Hence, the addition of very efficient IN in the right concen-
54 tration may result in fewer, larger ice crystals. The heterogeneously formed ice crystals
55 would deplete water vapor as they grow, and prevent the supersaturations required for the
56 onset of homogeneous freezing. Larger ice crystals would reduce cirrus optical thickness
57 and shorten cloud lifetimes through increased ice crystal sedimentation velocities. Both

58 mechanisms would yield a smaller greenhouse effect (Figure 1). *Mitchell and Finnegan*
59 [2009] proposed that the seeding material could be injected at cirrus levels by commer-
60 cial aircraft. A background concentration of seeding material would build up, and cirrus
61 clouds would form in an environment sufficiently enriched in IN for homogeneous freezing
62 to be suppressed. Here we have tested the effect of background concentrations of seeding
63 IN spanning several orders of magnitude in numerical simulations using a modified version
64 of the NCAR Community Atmosphere Model (CAM, version 5).

2. Modeling tool and experimental setup

65 CAM5 was for the purpose of this study run at a horizontal resolution of 1.9°latitude
66 and 2.5°longitude, with 30 vertical levels, a finite volume dynamical core and a timestep of
67 20 minutes. All simulations were conducted with climatological sea surface temperatures
68 corresponding to the year 2000. CAM5's predecessor is described in *Gent et al.* [2011],
69 but its cloud microphysics has since been updated [*Gettelman et al.*, 2008, 2010] and it
70 also has a recently developed modal aerosol treatment [*Liu et al.*, 2012]. The aerosol size
71 distribution can now be represented by either three or seven lognormal modes (MAM3
72 and MAM7, respectively). The treatment of cirrus cloud microphysics in CAM5 has also
73 been significantly improved relative to earlier versions [*Liu et al.*, 2007], but was in this
74 study partly replaced by an alternative and more flexible cirrus scheme, developed by
75 *Barahona and Nenes* [2008, 2009]. The scheme is based on an analytical solution of the
76 governing equations of a cooling air parcel. It explicitly accounts for the effect of cloud
77 formation conditions and aerosol properties on the cirrus ice crystal concentration. Com-
78 petition between homogeneous and heterogeneous ice nucleation, hence the influence of

ice nuclei on ice crystal concentration, is also accounted for. Heterogeneous ice nucleation is described through a generalized ice nucleation spectrum, which can have any functional form, providing flexibility in describing ice nucleation on different IN. While the CAM5 cirrus scheme has already been carefully validated, particularly in terms of global cloud and radiation fields [Liu *et al.*, 2012; Gettelman *et al.*, 2008, 2010], we show in Table 1 global averages of some key cloud and radiation fields for the standard CAM5, the modified CAM5 used in this study, as well as from observations. Both model versions were run with homogeneous nucleation only, for temperatures below -38°C . The introduction of the new cirrus scheme does not dramatically change the cloud and radiation fields. However, it does produce more ice crystals at cirrus levels, which leads to optically thicker and longer-lived cirrus clouds, hence the slightly larger ice and liquid water paths (the latter due to reduced accretion of liquid by falling ice crystals). Cirrus ice crystal concentrations lie in the range $10\text{-}1000\text{l}^{-1}$, with a global annual average at 200hPa of $\sim 400\text{l}^{-1}$. This is somewhat higher than the concentrations reported for example from the recent SPARTICUS campaign [Mitchell *et al.*, 2011], but values are very sensitive to the treatment of subgrid-scale vertical velocity (see Section 3). Here, we made the assumption that under unseeded conditions, cirrus clouds form solely through homogeneous ice nucleation. The concentration of solution droplets that could potentially nucleate homogeneously corresponds to the predicted number concentration of particles in the Aitken mode of the MAM3 aerosol module. Based on the number concentration and size of the solution droplets, as well as temperature and vertical velocity, the homogeneous nucleation rate was calculated [Barahona and Nenes, 2008]. The seeding IN were all conservatively

101 assumed to activate and nucleate ice at a supersaturation of 10% [Mitchell *et al.*, 2011].
 102 We have carried out 20 model simulations, each 10 years of length after a spin-up of 3
 103 months, in which the concentration of seeding IN in the UT (IN_s) were varied from 0 to
 104 1500 l^{-1} . Note that in the following, cirrus clouds refer to all clouds forming in the UT,
 105 which will here correspond to the part of the atmosphere with temperatures lower than
 106 -35°C

3. Results

107 The delivery method, dispersion and atmospheric fate of the seeding IN are beyond
 108 the scope of the present study. Here, we focus on the effect of seeding on cirrus cloud
 109 properties and Earth's energy budget, under the assumption that there exists some means
 110 to build up uniform background concentrations of seeding IN in the UT. Figure 2a shows
 111 simulated global annual mean vertically integrated ice amount (ice water path, IWP) and
 112 high cloud coverage (CC_{HGH}), as a function of IN_s . For low IN_s concentrations ($< 5 \text{ l}^{-1}$),
 113 IWP and CC_{HGH} remain very similar to their values under pure homogeneous freezing,
 114 i.e. $IN_s=0 \text{ l}^{-1}$, our reference case (REF). However, for IN_s in the range 5 l^{-1} - 100 l^{-1} , both
 115 are suppressed and ice crystals are 10-20% larger than in the case of pure homogeneous
 116 freezing (Fig. 2b). In this IN_s range we also observed a small reduction in liquid water
 117 path, due to increased accretion of liquid by falling ice crystals. Finally, for $IN_s > 100$
 118 l^{-1} , seeding leads to the opposite effect; smaller ice crystals and the consequent increase
 119 in IWP and CC_{HGH} . From Figure 2, three distinct regimes can be identified: (1) *The*
 120 *sub-optimal seeding regime*: IN_s is insufficient for suppression of homogeneous nucleation,
 121 and the cirrus clouds remain unaffected by the seeding, (2) *The optimal seeding regime*:

122 homogeneous nucleation is suppressed, and IN_s is low enough to reduce ice crystal concen-
123 tration and increase crystal size, with associated reductions in cirrus cloud amount and
124 coverage, and (3) *The over-seeding regime*: homogeneous nucleation is suppressed, but
125 more ice crystals nucleate on seeds than would otherwise have nucleated homogeneously
126 in the unseeded case. Table 2 gives approximate IN_s intervals for these three regimes in
127 our control model set-up (CTL). As a consequence of the increase in ice crystal sizes and
128 decrease in cirrus cloud amount in the optimal seeding regime, cirrus clouds become op-
129 tically thinner, as illustrated by the reduction in longwave cloud forcing (LWCF), shown
130 in Figure 2c. The reduced LWCF allows for more outgoing longwave radiation at the
131 top-of-the atmosphere (TOA), corresponding to a negative radiative forcing (i.e. cool-
132 ing) of about 7 Wm^{-2} . This cooling is partly compensated for by a reduction in cirrus
133 cloud albedo and hence the shortwave cloud forcing (SWCF), such that the maximum
134 reduction in the net cloud forcing (NCF) amounts to 2.0 Wm^{-2} . While changes in the net
135 shortwave flux at the TOA are very similar to the changes in SWCF, the reduction in UT
136 water vapor in response to the seeding increases the outgoing longwave radiation further
137 by up to 0.5 Wm^{-2} , and hence amplifies the cooling. Hence, cirrus cloud seeding could
138 potentially eliminate a forcing equivalent to that which has been causing climate change
139 to date. However, this would require seeding IN concentrations finely tuned to lie exactly
140 in the optimal IN_s window. While the main perceived risk of under-seeding is a costly,
141 wasted effort, over-seeding could actually lead to the opposite of the desired effect. This
142 is illustrated in Figure 2; IN_s concentrations larger than 100 l^{-1} would lead to an increase
143 in IWP and a decrease in ice crystal sizes relative to the unseeded atmosphere, and hence

144 a warming rather than a cooling. Based on Figure 2, we have approximated the optimal
 145 IN_s , $IN_{s,o}$, to $15l^{-1}$, and have displayed anomalies in several cirrus cloud properties rela-
 146 tive to REF for $IN_{s,o}$ in Figure 3. Evident is the strong reduction in ice crystal number
 147 concentrations in the UT (Fig. 3a), which allows individual ice crystals to grow larger
 148 via vapour deposition (Fig. 3b). The larger ice crystals in turn lead to reduced cloud ice
 149 (Fig. 3c) and cloud coverage (Fig. 3d), as a result of the faster sedimentation of the larger
 150 ice crystals. As expected, the strongest perturbations are found at mid-latitudes, where
 151 cirrus clouds form in situ, rather than in the tropics, where anvil cirrus are produced by
 152 convective outflow.

153 Several studies have indicated that the relative importance of homogeneous vs. heteroge-
 154 neous ice nucleation is very sensitive to the vertical velocity at the cloud-scale [*Kaercher*
 155 *and Lohmann, 2003; DeMott et al., 1997*]. CESM parameterizes this subgrid-scale updraft
 156 velocity as a single value for each model grid box, proportional to the square root of the
 157 turbulent kinetic energy (TKE), $W_c = \sqrt{\frac{2}{3}TKE}$. We have tested the robustness of our
 158 results to increased/decreased vertical velocities, by repeating the set of IN_s perturba-
 159 tion simulations, but with $W_c = W_{c,HGH} = \sqrt{\frac{8}{3}TKE}$ and $W_c = W_{c,LOW} = \sqrt{\frac{1}{6}TKE}$,
 160 respectively. Figure 4 shows the change in NCF (relative to REF) as a function of IN_s for
 161 simulations with W_c , $W_{c,HGH}$ and $W_{c,LOW}$. Evident from Table 1 is a shift in the optimal
 162 IN_s interval towards lower (higher) values when W_c is decreased (increased). The mag-
 163 nitude of the cooling is also affected, and becomes smaller (larger) when W_c is decreased
 164 (increased). Higher vertical velocities lead to higher homogeneous nucleation rates, and
 165 hence a stronger perturbation when homogeneous nucleation is suppressed. Higher verti-

166 cal velocities also require higher IN_s concentrations in order for homogeneous nucleation
167 to be suppressed. While previous studies of the effect of anthropogenic IN on cirrus have
168 reported a sensitivity to the concentration of solution droplets available for homogeneous
169 nucleation[*Penner et al.*, 2009], we found minor changes in a simulation reducing the
170 concentration of solution droplets available by 50%.

4. Discussion and outlook

171 Further investigations of the viability of cirrus seeding as a means of stabilizing Earth's
172 climate will require simulations of the atmospheric lifetimes of seeding IN, from the point of
173 emission, through potential ice nucleation, and subsequent sedimentation and deposition
174 on Earth's surface. Laboratory investigations of ice nucleation on BiI3 are also required
175 to shed further light on the geoengineering process investigated here. The present study
176 has demonstrated that successful cirrus cloud seeding requires seeding IN concentrations
177 that lie in a relatively narrow optimal interval. The bounds of this interval are set by
178 the vertical velocities in the UT, for which only sparse and sporadic measurements exist.
179 A premature implementation of cirrus seeding before knowledge of vertical velocities at
180 cirrus levels is improved could accelerate global warming as opposed to prevent it.

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Table 1. Simulated global and annual mean cloud cover (CC), ice water path (IWP), liquid water path (LWP) and net cloud forcing (NCF) from the standard and modified CAM5.1 (CAM5.1-HOM and CAM5.1-BN09, respectively) as well as from satellite observations (OBS). Observations are taken from a combination of CloudSat and CALIPSO retrievals (CC, IWP and LWP), and from ERBE and CERES (NCF).

	CC (%)	IWP (gm^{-2})	LWP (gm^{-2})	NCF (Wm^{-2})
CAM5.1-HOM	64.3	17.8	44.2	-27.6
CAM5.1-BN0	68.8	21.9	47.1	-26.3
OBS	71	20 to 70	30 to 50	-17.2 to -23.8

Table 2. Approximate sub-optimal, optimal, and over-seeding INs concentrations for the CTL, $W_{c,LOW}$ and $W_{c,HGH}$ sets of simulations.

CASE	Sub-optimal IN_s	Optimal IN_s	Over-seeding IN_s
CTL	$<5\text{l}^{-1}$	$5\text{-}100\text{l}^{-1}$	$>100\text{l}^{-1}$
$W_{c,LOW}$	$<1\text{l}^{-1}$	$1\text{-}25\text{l}^{-1}$	$>25\text{l}^{-1}$
$W_{c,HGH}$	$<20\text{l}^{-1}$	$20\text{-}200\text{l}^{-1}$	$>200\text{l}^{-1}$

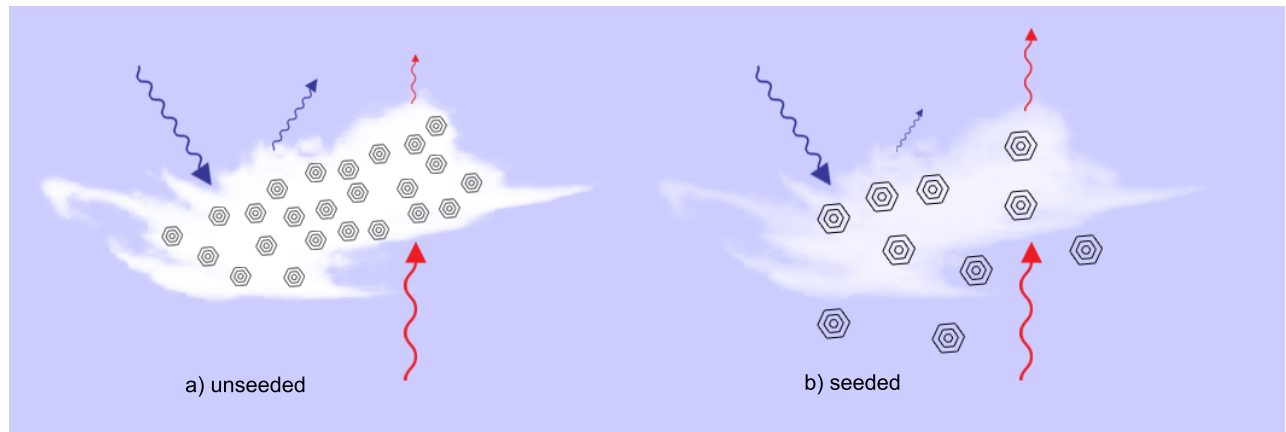


Figure 1. Conceptual schematic of changes in cirrus cloud properties in response to seeding. Red arrows represent longwave (LW) radiation and blue arrows represent shortwave (SW) radiation. The seeded cirrus clouds on average reflect slightly less SW radiation back to Space, but also allow more LW radiation to escape to Space, and the latter effect dominates.

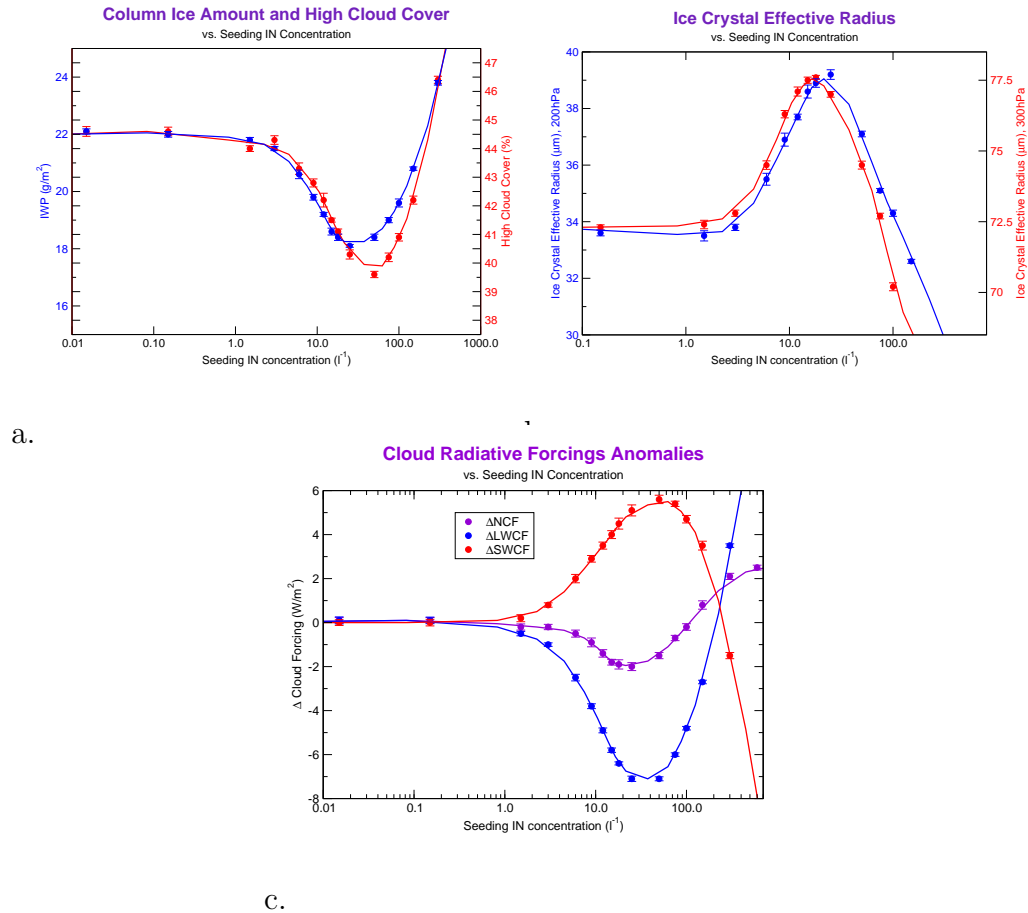


Figure 2. CESM simulations of macro-physical and radiative properties of high clouds as a function of IN_s . Each circle corresponds to an individual 2-year CESM simulation. a) High cloud amount (i.e. cloud cover integrated from 400hPa to 50hPa) and vertically integrated ice amount (Ice Water Path, IWP), b) Ice crystal effective radius at 300hPa (red solid line) and 200 hPa (blue solid line), and c) changes in longwave-, shortwave- and net cloud forcing (SWCF, LWCF and NCF, respectively) at the top-of-the-atmosphere (TOA), relative to REF. Solid lines represent moving averages. Error bars represent one standard deviation, calculated from annual averages.

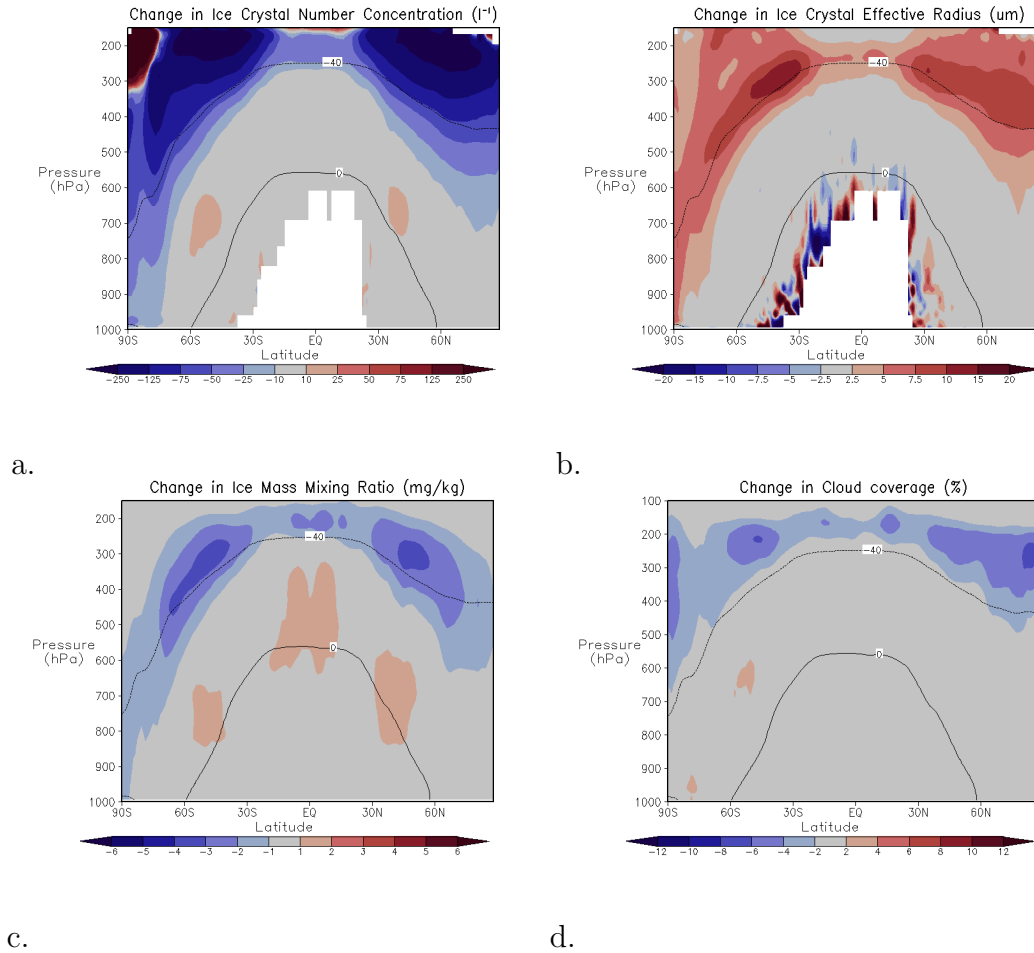


Figure 3. Simulated changes in zonal and annual mean cloud properties induced by a seeding IN concentration of $15l^{-1}$ (relative to REF): a) in-cloud ice crystal number concentration, b) ice crystal effective radius, c) ice mass mixing ratio and d) cloud coverage. All plots are based on 10-year model simulations.

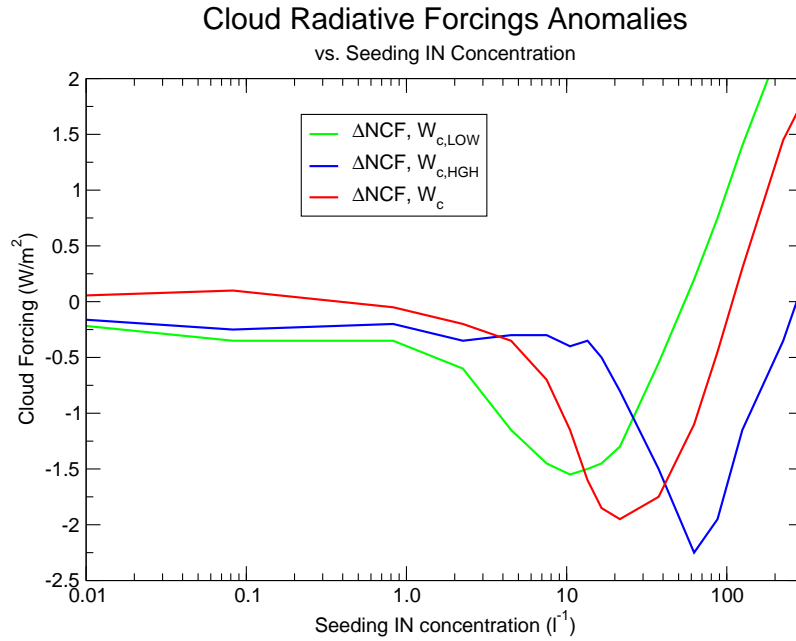


Figure 4. Change in the net cloud forcing (NCF) as a function of IN_s at the TOA relative to REF for default, doubled and halved subgrid-scale vertical velocity (W_c , $W_{c,HGH}$, $W_{c,LOW}$, respectively).